

H-Toeplitz算子的代数性质

梁金金

浙江师范大学, 数学与计算机科学学院, 浙江 金华

收稿日期: 2021年11月27日; 录用日期: 2021年12月23日; 发布日期: 2021年12月30日

摘要

本文主要研究 Bergman 空间上 H-Toeplitz 算子的代数性质。第一章介绍了相关的研究背景、基本概念及一些主要结果。第二章给出了本文主要结果的证明, 证明了拟齐次符号 H-Toeplitz 算子的复对称性。

关键词

Bergman空间, H-Toeplitz算子, 复对称性

Algebraic Properties of H-Toeplitz Operator

Jinjin Liang

College of Mathematics and Computer Science, Zhejiang Normal University, Jinhua Zhejiang

Received: Nov. 27th, 2021; accepted: Dec. 23rd, 2021; published: Dec. 30th, 2021

Abstract

In this paper, we mainly study the algebraic properties of H-toeplitz operators on Bergman Spaces. In Chapter 1, we introduce the related research background, basic

concepts and some main results. In Chapter 2, the proof of the main results of this paper is given, and the complex symmetry of quasi homogeneous signed H-toeplitz operators is proved.

Keywords

Bergman Space, H-Toeplitz Operator, Complex Symmetry

Copyright © 2021 by author(s) and Hans Publishers Inc.

This work is licensed under the Creative Commons Attribution International License (CC BY 4.0).

<http://creativecommons.org/licenses/by/4.0/>



Open Access

1. 介绍

记 D 为复平面 C 内的单位开圆盘 $dA(z) = \frac{1}{\pi} r dr d\theta$ 为 D 上规范化的 Lebesgue 面积测度. 记 $H(D)$ 为 D 上所有解析函数, 记 Bergman 空间 $L_a^2(D)$ 为 D 上所有勒贝格平方可积解析函数所成的 Hilbert 空间, 其内积定义为

$$\langle f, g \rangle = \int_D f(z) \overline{g(z)} dA(z), \quad f, g \in L^2(D).$$

空间 $L_a^2(D)$ 是 Hilbert 空间 $L^2(D, dA)$ 的闭子空间. 对于非负整数 n , 设 $e_n(z) = \sqrt{n+1} z^n$. 对所有 $z \in D$. 集合 $\{e_n\}_{n \geq 0}$ 构成 $L^2(D)$ 的一组标准正交基 [1]. 对于 $z, w \in D$, Bergman 空间中的再生核称为 Bergman 核, 定义为

$$K_z(w) = K(z, w) = \frac{1}{(1 - \bar{z}w)^2}.$$

记 $P_{L_a^2} : L^2(D, dA) \rightarrow L_a^2(D)$, 表示空间 $L^2(D, dA)$ 到空间 $L_a^2(D)$ 的正交投影, 称为 Bergman 投影. 对所有 $f \in L^2(D, dA)$, 由再生核性质:

$$P_{L_a^2}(f)(z) = \langle P_{L_a^2} f, K_z \rangle = \langle f, K_z \rangle = \int_D \frac{f(w)}{(1 - z\bar{w})^2} dA(w).$$

见 [2]. 对 $\phi \in L^\infty(D)$, 乘法算子 M_ϕ 定义为 $M_\phi(f) = \phi f$. 符号为 $\phi \in L^\infty(D)$ 的 Toeplitz 算子 $T_\phi : L_a^2(D) \rightarrow L_a^2(D)$ 由下式定义

$$T_\phi(f) = P_{L_a^2}(\phi f) = \int_D \frac{\phi(z) f(w)}{(1 - \bar{z}w)^2} dA(w), \quad f \in L^2(D).$$

对 $\phi \in L^\infty(D)$, Hankel 算子 $H_\phi : L_a^2(D) \rightarrow L_a^2(D)$ 定义为

$$H_\phi(f) = P_{L_a^2} M_\phi J(f), \forall f \in L_a^2(D),$$

其中算子 $J : L_a^2(D) \rightarrow \overline{L_a^2(D)}$ 由 $J(e_n(z)) = \overline{e_{n+1}(z)}$ 给出, 其中 n 为非负整数. T_ϕ 和 H_ϕ 是 $L_a^2(D)$ 上的有界线性算子.

引理 1.1 在 Bergman 空间 $L_a^2(D)$ 中对于非负整数 s 和 t , 以下成立, 见 [3]:

$$\langle z^s, z^t \rangle = \begin{cases} \frac{1}{s+1}, & s=t; \\ 0, & s \neq t. \end{cases}$$

$$P_{L_a^2}(\overline{z^t} z^s) = \begin{cases} \frac{s-t+1}{s+1} z^{s-t}, & s \geq t; \\ 0, & s < t. \end{cases}$$

为了定义 $L_a^2(D)$ 上的 H-Toeplitz 算子的概念, 首先定义线性算子 $K : L_a^2(D) \rightarrow L_{harm}^2(D)$,

$$K(e_{2n}) = e_n, \quad K(e_{2n+1}) = \overline{e_{n+1}}, \quad n = 0, 1, 2, \dots$$

易见算子 K 在 $L_a^2(D)$ 上有界. 此外, $\|K\| = 1$. 算子 K 的伴随由下式给出: 对 $\forall n \geq 0$,

$$K^*(e_n) = e_{2n}, \quad K^*(\overline{e_{n+1}}) = e_{2n+1}.$$

对 $\phi \in L^\infty(D)$, H-Toeplitz 算子定义为 $B_\phi : L_a^2(D) \rightarrow L_a^2(D)$,

$$B_\phi(f) = P_{L_a^2} M_\phi K(f), f \in L_a^2(D),$$

对于每个非负整数 n , 我们有

$$B_\phi(e_{2n}) = P_{L_a^2} M_\phi K(e_{2n}) = P_{L_a^2} M_\phi(e_n) = T_\phi(e_n)$$

和

$$B_\phi(e_{2n+1}) = P_{L_a^2} M_\phi K(e_{2n+1}) = P_{L_a^2} M_\phi(e_{n+1}) = P_{L_a^2} M_\phi J(e_n) = H_\phi(e_n).$$

故 H-Toeplitz 算子与 Toeplitz 算子与 Hankel 算子密切相关. 见 [4] [5].

2. 主要结果及证明

关于 $L_a^2(D)$ 的标准正交基 $\{e_n\}_{n \geq 0}$, T_ϕ 矩阵的 $(m, n)^{th}$, 由下式给出:

$$\begin{aligned}
\langle T_\phi(e_n), e_m \rangle &= \langle P_{L_a^2}(\phi e_n), e_m \rangle \\
&= \sqrt{n+1}\sqrt{m+1} \langle (\sum_{i=0}^{\infty} a_i z^i + \sum_{j=1}^{\infty} b_j \bar{z}^j) z^n, z^m \rangle \\
&= \sqrt{n+1}\sqrt{m+1} (\sum_{i=0}^{\infty} a_i \langle z^{i+n}, z^m \rangle + \sum_{j=1}^{\infty} \langle b_j z^n, z^{m+j} \rangle).
\end{aligned}$$

分以下两种情形:

情形(1): 如果 $m \geq n$, 我们有:

$$\begin{aligned}
\langle T_\phi(e_n), e_m \rangle &= \sqrt{n+1}\sqrt{m+1} \sum_{i=0}^{\infty} a_i \langle z^{i+n}, z^m \rangle \\
&= \sqrt{n+1}\sqrt{m+1} \frac{1}{m+1} a_{m-n} \\
&= \sqrt{\frac{n+1}{m+1}} a_{m-n}.
\end{aligned}$$

情形(2): 如果 $m < n$, 我们有:

$$\begin{aligned}
\langle T_\phi(e_n), e_m \rangle &= \sqrt{n+1}\sqrt{m+1} \sum_{j=1}^{\infty} \langle b_j z^n, z^{m+j} \rangle \\
&= \sqrt{n+1}\sqrt{m+1} \frac{1}{n+1} b_{n-m} \\
&= \sqrt{\frac{m+1}{n+1}} b_{n-m}.
\end{aligned}$$

因此, 我们可以得到:

$$\langle T_\phi(e_n), e_m \rangle = \begin{cases} \sqrt{\frac{n+1}{m+1}} a_{m-n}, & m \geq n; \\ \sqrt{\frac{m+1}{n+1}} b_{n-m}, & m \leq n. \end{cases}$$

其中 m 和 n 是非负整数, H_ϕ 的矩阵 $(m, n)^{th}$ 相对于 $L_a^2(D)$ 的标准正交基 $\{e_n\}_{n \geq 0}$ 由下式给出:

$$\begin{aligned}
\langle H_\phi(e_n), e_m \rangle &= \langle P_{L_a^2} M_\phi J(e_n), e_m \rangle \\
&= \sqrt{n+2}\sqrt{m+1} \langle (\sum_{i=0}^{\infty} a_i z^i + \sum_{j=1}^{\infty} b_j \bar{z}^j) \bar{z}^{n+1}, z^m \rangle \\
&= \sqrt{n+2}\sqrt{m+1} (\sum_{i=0}^{\infty} a_i \langle z^i, z^{m+n+1} \rangle + \sum_{j=1}^{\infty} b_j \langle \bar{z}^j, z^{m+n+1} \rangle) \\
&= \frac{\sqrt{m+1}\sqrt{n+2}}{m+n+2} a_{m+n+1}.
\end{aligned}$$

同理,对于非负整数 m 和 n , 有

$$\langle B_\phi(e_{2n}), e_m \rangle = \langle T_\phi(e_n), e_m \rangle \begin{cases} \sqrt{\frac{n+1}{m+1}} a_{m-n}, & m \geq n; \\ \sqrt{\frac{m+1}{n+1}} b_{n-m}, & m \leq n. \end{cases}$$

和

$$\langle B_\phi(e_{2n+1}), e_m \rangle = \langle H_\phi(e_n), e_m \rangle = \frac{\sqrt{m+1}\sqrt{n+2}}{m+n+2} a_{m+n+1},$$

其中 m 和 n 是非负整数, 因此, B_ϕ 关于 $L^2_a(D)$ 的标准正交基 $\{e_n\}_{n \geq 0}$ 的矩阵表示由下式给出:

$$B_\phi = \begin{bmatrix} a_0 & \frac{1}{\sqrt{2}}a_1 & \frac{1}{\sqrt{2}}b_1 & \frac{1}{\sqrt{3}}a_2 & \frac{1}{\sqrt{3}}b_2 & \frac{1}{2}a_3 & \frac{1}{2}b_3 & \cdots \\ \frac{1}{\sqrt{2}}a_1 & \frac{1}{\sqrt{3}}a_2 & a_0 & \frac{\sqrt{6}}{4}a_3 & \sqrt{\frac{2}{3}}b_1 & \frac{2\sqrt{2}}{5}a_4 & \frac{1}{\sqrt{2}}b_2 & \cdots \\ \frac{1}{\sqrt{3}}a_2 & \frac{\sqrt{6}}{4}a_3 & \sqrt{\frac{2}{3}}a_1 & \frac{3}{5}a_4 & a_0 & \frac{1}{\sqrt{3}}a_5 & \frac{\sqrt{3}}{2}b_1 & \cdots \\ \frac{1}{2}a_3 & \frac{2\sqrt{2}}{5}a_4 & \frac{1}{\sqrt{2}}a_2 & \frac{1}{\sqrt{3}}a_5 & 3\frac{\sqrt{3}}{2}a_1 & \frac{4}{7}a_6 & a_0 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}.$$

它的伴随矩阵如下:

$$B_\phi^* = \begin{bmatrix} \bar{a}_0 & \frac{1}{\sqrt{2}}\bar{a}_1 & \frac{1}{\sqrt{3}}\bar{a}_2 & \frac{1}{2}\bar{a}_3 & \cdots \\ \frac{1}{\sqrt{2}}\bar{a}_1 & \frac{2}{3}\bar{a}_2 & \frac{\sqrt{6}}{4}\bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}}\bar{a}_4 & \cdots \\ \frac{1}{\sqrt{2}}\bar{b}_1 & \bar{a}_0 & \frac{\sqrt{3}}{2}\bar{a}_1 & \frac{1}{\sqrt{2}}\bar{a}_2 & \cdots \\ \frac{1}{\sqrt{3}}\bar{a}_2 & \frac{\sqrt{6}}{4}\bar{a}_3 & \frac{3}{5}\bar{a}_4 & \frac{1}{\sqrt{3}}\bar{a}_5 & \cdots \\ \frac{1}{\sqrt{3}}\bar{b}_2 & \frac{\sqrt{3}}{2}\bar{b}_1 & \bar{a}_0 & \frac{\sqrt{3}}{2}\bar{a}_1 & \cdots \\ \frac{1}{2}\bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}}\bar{a}_4 & \frac{1}{\sqrt{3}}\bar{a}_5 & \frac{4}{7}\bar{a}_6 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}.$$

空间 $L^2(D, dA)$ 上的复共轭算子 $C : L^2(D) \rightarrow L^2(D)$, 对于 $f, g \in L^2(D, dA)$, 满足 $C^2 = I$ 和 $\langle Cf, Cg \rangle = \langle g, f \rangle$ 都成立. [6] 证明了 $L^2(D, dA)$ 存在一个正交基 $\{e_n\}_{n=0}^\infty$, 使得 $Ce_i = e_i$, $i = 0, 1, \dots$. 若 $L^2(D, dA)$ 上的一个有界线性算子 T , 如果存在复共轭 C , 使得 $CT = T^*C$, 则称 T 是复对称的. 见 [7].

定理 2.1 令 $\phi \in L^\infty(D)$ 为调和符号. 那么 H -Toeplitz 算子 B_ϕ 在 $L^2(D, dA)$ 上关于 C 为复对称算子当且仅当 $\phi = 0$.

证明. 设 $\phi(z) = \sum_{i=0}^\infty a_i z^i + \sum_{j=0}^\infty b_j \bar{z}^j \in L^\infty(D)$ 为 D 上的调和函数, 并设 B_ϕ 是复对称算子. 这

意味着 $CB_\phi C = B_\phi^*$. 对于每个非负整数 n , 通过引理1.1 我们可以计算出:

$$\begin{aligned}
 \langle CB_\phi C e_{2n}, e_m \rangle &= \langle C e_m, B_\phi C e_{2n} \rangle \\
 &= \langle e_m, B_\phi e_{2n} \rangle = \langle e_m, P_{L_a^2} M_\phi K(e_{2n}) \rangle \\
 &= \langle e_m, P_{L_a^2} M_\phi(e_n) \rangle = \langle e_m, M_\phi(e_n) \rangle \\
 &= \langle \sqrt{m+1} z^m, (\sum_{i=0}^{\infty} a_i z^i + \sum_{j=0}^{\infty} b_j \bar{z}^j) \sqrt{n+1} z^n \rangle \\
 &= \sqrt{m+1} \sqrt{n+1} \langle z^m, \sum_{i=0}^{\infty} a_i z^i z^n \rangle + \sqrt{m+1} \sqrt{n+1} \langle z^m, \sum_{j=0}^{\infty} b_j \bar{z}^j z^n \rangle.
 \end{aligned}$$

因此, 出现了以下两种情形:

情形(1): 如果 $m \geq n$, 那么

$$\begin{aligned}
 \langle CB_\phi C e_{2n}, e_m \rangle &= \langle \sqrt{m+1} \sqrt{n+1} z^m, \sum_{i=0}^{\infty} a_i z^i z^n \rangle \\
 &= \sqrt{m+1} \sqrt{n+1} \sum_{i=0}^{\infty} \bar{a}_i \langle z^m, z^{i+n} \rangle \\
 &= \sqrt{m+1} \sqrt{n+1} \frac{1}{m+1} \bar{a}_{m-n} \\
 &= \sqrt{\frac{n+1}{m+1}} \bar{a}_{m-n}.
 \end{aligned}$$

情形(2): 如果 $m < n$, 那么

$$\begin{aligned}
 \langle CB_\phi C e_{2n}, e_m \rangle &= \sqrt{m+1} \sqrt{n+1} \langle z^m, \sum_{j=0}^{\infty} b_j \bar{z}^j z^n \rangle \\
 &= \sqrt{m+1} \sqrt{n+1} \sum_{j=0}^{\infty} \bar{b}_j \langle z^{m+j}, z^n \rangle \\
 &= \sqrt{m+1} \sqrt{n+1} \frac{1}{n+1} \bar{b}_{n-m} \\
 &= \sqrt{\frac{m+1}{n+1}} \bar{b}_{n-m}.
 \end{aligned}$$

$$\begin{aligned}
\langle CB_\phi C e_{2n+1}, e_m \rangle &= \langle C e_m, B_\phi C e_{2n+1} \rangle \\
&= \langle e_m, B_\phi \overline{e_{2n+1}} \rangle = \langle e_m, P_{L_a^2} M_\phi K(e_{2n+1}) \rangle \\
&= \langle e_m, P_{L_a^2} M_\phi \overline{e_{n+1}} \rangle = \langle e_m, M_\phi \overline{e_{n+1}} \rangle = \langle e_m, \phi \overline{e_{n+1}} \rangle \\
&= \langle \sqrt{m+1} z^m, (\sum_{i=0}^{\infty} a_i z^i + \sum_{j=0}^{\infty} b_j \bar{z}^j) \sqrt{n+2} \bar{z}^{n+1} \rangle \\
&= \sqrt{m+1} \sqrt{n+1} \langle z^m, \sum_{i=0}^{\infty} a_i z^i \bar{z}^{n+1} \rangle + \langle z^m, \sum_{j=0}^{\infty} b_j \bar{z}^j \bar{z}^{n+1} \rangle \\
&= \sqrt{m+1} \sqrt{n+1} \sum_{i=0}^{\infty} \bar{a}_i \langle z^m, z^i \bar{z}^{n+1} \rangle \\
&= \frac{\sqrt{m+1} \sqrt{n+2}}{n+m+2} a_{m+n+1}.
\end{aligned}$$

其中 m 和 n 是非负整数. 因此, 可得到 $CB_\phi C$ 的矩阵形式如下:

$$CB_\phi C = \begin{bmatrix} \bar{a}_0 & \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{1}{\sqrt{2}} \bar{b}_1 & \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{1}{\sqrt{3}} \bar{b}_2 & \frac{1}{2} \bar{a}_3 & \frac{1}{2} \bar{b}_3 \cdots \\ \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{2}{3} \bar{a}_2 & \bar{a}_0 & \frac{\sqrt{6}}{4} \bar{a}_3 & \frac{1}{\sqrt{2}} \bar{b}_1 & \frac{2\sqrt{2}}{\sqrt{5}} \bar{a}_4 & \frac{1}{\sqrt{2}} \bar{b}_2 \cdots \\ \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{\sqrt{6}}{4} \bar{a}_3 & \sqrt{\frac{2}{3}} \bar{a}_1 & \frac{3}{5} \bar{a}_4 & \bar{a}_0 & \frac{1}{\sqrt{3}} \bar{a}_5 & \frac{\sqrt{3}}{2} \bar{b}_1 \cdots \\ \frac{1}{2} \bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}} \bar{a}_4 & \frac{1}{\sqrt{2}} \bar{a}_2 & \frac{1}{\sqrt{3}} \bar{a}_5 & \frac{\sqrt{3}}{2} \bar{a}_1 & \frac{4}{7} \bar{a}_6 & \bar{a}_0 \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \end{bmatrix}.$$

$$B_\phi^* = \begin{bmatrix} \bar{a}_0 & \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{1}{2} \bar{a}_3 & \cdots \\ \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{2}{3} \bar{a}_2 & \frac{\sqrt{6}}{4} \bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}} \bar{a}_4 & \cdots \\ \frac{1}{\sqrt{2}} \bar{b}_1 & \bar{a}_0 & \frac{\sqrt{3}}{2} \bar{a}_1 & \frac{1}{\sqrt{2}} \bar{a}_2 & \cdots \\ \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{\sqrt{6}}{4} \bar{a}_3 & \frac{3}{5} \bar{a}_0 & \frac{\sqrt{3}}{2} \bar{a}_1 & \cdots \\ \frac{1}{2} \bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}} \bar{a}_4 & \frac{1}{\sqrt{3}} \bar{a}_5 & \frac{4}{7} \bar{a}_6 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots \end{bmatrix}.$$

所以, 如果 $CB_\phi C = B_\phi^*$, 那么

$$\begin{bmatrix} \bar{a}_0 & \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{1}{\sqrt{2}} \bar{b}_1 & \frac{1}{\sqrt{3}} \bar{a}_2 \cdots \\ \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{2}{3} \bar{a}_2 & \bar{a}_0 & \frac{\sqrt{6}}{4} \bar{a}_3 \cdots \\ \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{\sqrt{6}}{4} \bar{a}_3 & \sqrt{\frac{2}{3}} \bar{a}_1 & \frac{3}{5} \bar{a}_4 \cdots \\ \frac{1}{2} \bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}} \bar{a}_4 & \frac{1}{\sqrt{2}} \bar{a}_2 & \frac{1}{\sqrt{3}} \bar{a}_5 \cdots \\ \vdots & \vdots & \vdots & \vdots \end{bmatrix} = \begin{bmatrix} \bar{a}_0 & \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{1}{2} \bar{a}_3 \cdots \\ \frac{1}{\sqrt{2}} \bar{a}_1 & \frac{2}{3} \bar{a}_2 & \frac{\sqrt{6}}{4} \bar{a}_3 & \frac{2\sqrt{2}}{\sqrt{5}} \bar{a}_4 \cdots \\ \frac{1}{\sqrt{2}} \bar{b}_1 & \bar{a}_0 & \frac{\sqrt{3}}{2} \bar{a}_1 & \frac{1}{\sqrt{2}} \bar{a}_2 \cdots \\ \frac{1}{\sqrt{3}} \bar{a}_2 & \frac{\sqrt{6}}{4} \bar{a}_3 & \frac{3}{5} \bar{a}_4 & \frac{1}{\sqrt{3}} \bar{a}_5 \cdots \\ \vdots & \vdots & \vdots & \vdots \end{bmatrix}.$$

当 $n > m$ 时,

$$\begin{aligned}\sqrt{\frac{m+1}{n+1}}a_{n-m} &= \sqrt{\frac{m+1}{n+1}}b_{n-m}; \\ \sqrt{\frac{m+1}{n+1}}a_{n-m} &= \frac{\sqrt{m+1}\sqrt{m+2}}{m+n+2}a_{m+n+1}; \\ \sqrt{\frac{m+1}{n+1}}b_{n-m} &= \frac{\sqrt{m+1}\sqrt{n+2}}{m+n+2}a_{m+n+1}.\end{aligned}$$

当 $n \leq m$ 时,

$$\begin{aligned}\sqrt{\frac{n+1}{m+1}}a_{m-n} &= \sqrt{\frac{n+1}{m+1}}b_{m-n}; \\ \sqrt{\frac{n+1}{m+1}}a_{m-n} &= \frac{\sqrt{m+1}\sqrt{n+2}}{n+m+2}a_{m+n+1}; \\ \sqrt{\frac{n+1}{m+1}}b_{m-n} &= \frac{\sqrt{m+1}\sqrt{n+2}}{n+m+2}a_{m+n+1}.\end{aligned}$$

故对所有非负整数 m 和 n , 满足 $a_1 = a_2 = a_3 = \cdots = a_{n-m} = a_{m+n+1} \rightarrow \infty$. $a_1 = b_1 = a_2 = b_2 = a_3 = b_3 = \cdots = a_{n-m} = b_{n-m} = a_{m+n+1} \rightarrow \infty$, 当 $m, n \rightarrow \infty$. 矩阵中的每个元素都包含, 又因为在 $L^2(D)$ 中矩阵满足 $\sup \sum_{i=0}^{\infty} |a_i|^2 < \infty$, $\sup \sum_{j=0}^{\infty} |b_j|^2 < \infty$, 所以当 $m, n \rightarrow \infty$, $\lim_{i \rightarrow \infty} \sum_{i=0}^{\infty} |a_i|^2 = 0$, $\lim_{j \rightarrow \infty} \sum_{j=0}^{\infty} |b_j|^2 = 0$, 这意味着对于所有 i 和 j , $a_i = 0$ 且 $b_j = 0$. 因此我们得到了期望的结果 $\phi = 0$.

参考文献

- [1] Ahern, P. and Čučković, Ž. (2001) A Theorem of Brown-Halmos Type for Bergman Space Toeplitz Operators. *Journal of Functional Analysis*, **187**, 200-210. <https://doi.org/10.1006/jfan.2001.3811>
- [2] Brown, A. and Halmos, P.R. (1964) Algebraic Properties of Toeplitz Operators. *Journal für die Reine und Angewandte Mathematik*, **213**, 89-102. <https://doi.org/10.1515/crll.1964.213.89>
- [3] Louhichi, I. and Zakariasy, L. (2005) On Toeplitz Operators with Quasihomogeneous Symbols. *Archiv der Mathematik (Basel)*, **85**, 248-257. <https://doi.org/10.1007/s00013-005-1198-0>
- [4] Gupta, A. and Singh, S.K. (2021) H-Toeplitz Operators on Bergman Space. *Bulletin of the Korean Mathematical Society*, **58**, 327-347.
- [5] Lu, Y. and Zhang, B. (2011) Commuting Hankel Operator and Toeplitz Operator on the Bergman Space. *Chinese Annals of Mathematics, Series A*, **32**, 519-530.
- [6] Chen, Y., Lee, Y.J. and Zhao, Y. (2021) Complex Symmetry of Toeplitz Operators. Preprint.

- [7] Garcia, S.R. and Putinar, M. (2006) Complex Symmetric Operators and Applications. *Transactions of the American Mathematical Society*, **358**, 1285-1315.
<https://doi.org/10.1090/S0002-9947-05-03742-6>